



PANDA Electromagnetic Working Group Meeting
Ferrara, 15 – 16 October 2007



DRELL YAN AT PANDA

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Introduction

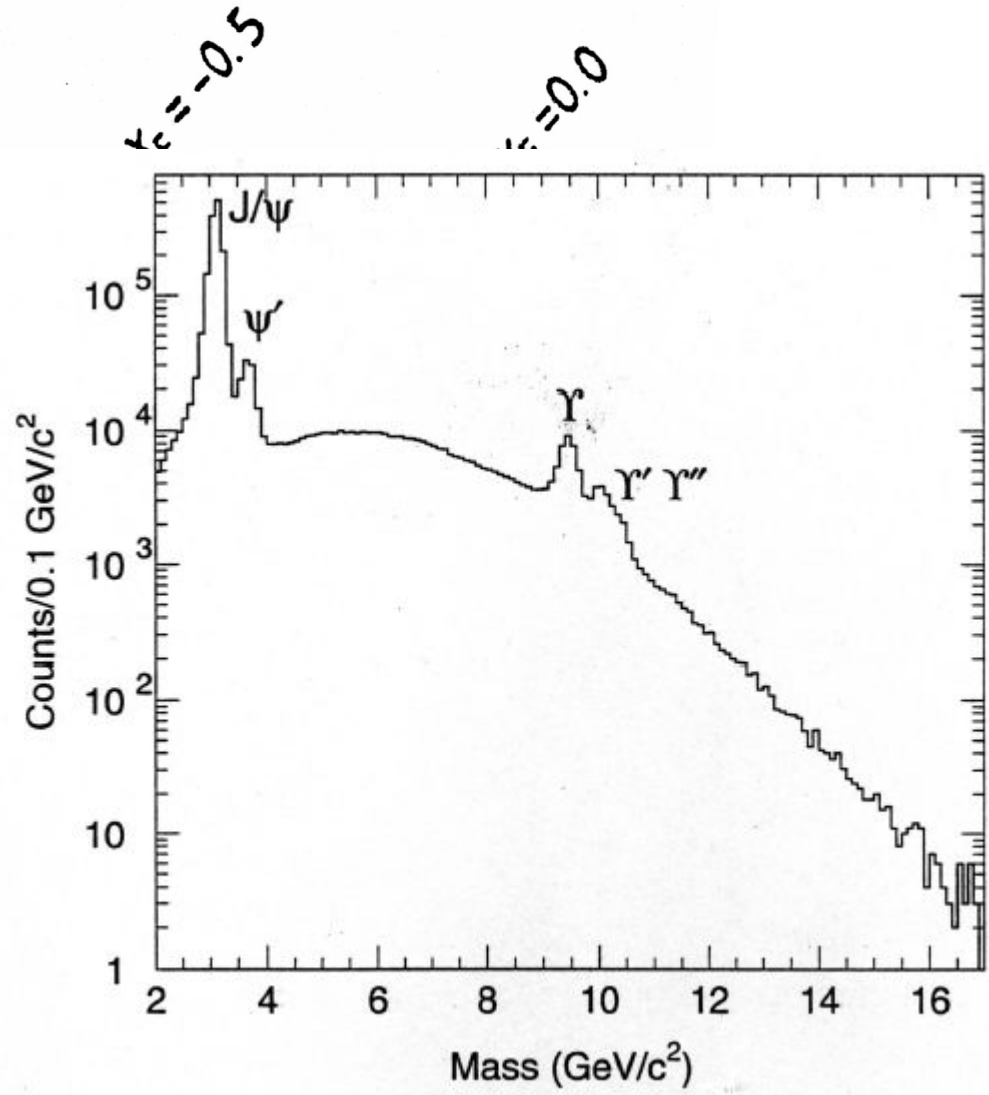
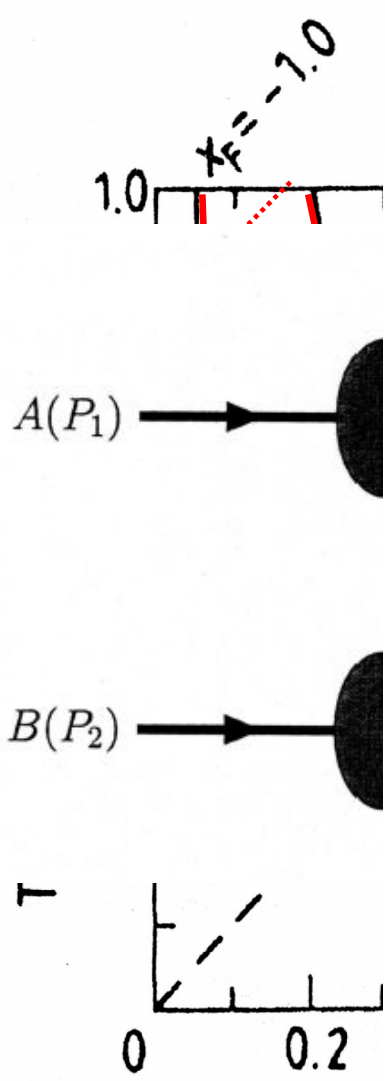


A complete description of nucleonic structure requires:

- quark and gluon distribution functions
- quark fragmentation functions

@ leading twist and @ NLO (k_T dependence)

Phase space for Drell-Yan processes



hyperbolae
kinematics

$$x_2 = \frac{M^2}{2P_2 \cdot q}$$

$$x_F = x_1 - x_2$$

$$x_1 x_2 = \frac{M^2}{s}$$

Beam antiquark / quark, x

Drell-Yan Asymmetries — $\bar{p}p \rightarrow \mu^+ \mu^- X$

Di-Lepton Rest Frame

$$\frac{1}{\sigma} \frac{d\sigma}{d\Omega} = \frac{3}{4\pi} \frac{1}{\lambda + 3} \left(1 + \lambda \cos^2 \theta + \mu \sin^2 \theta \cos \varphi + \frac{\nu}{2} \sin^2 \theta \cos 2\varphi \right)$$

NLO pQCD: $\lambda \sim 1, \mu \sim 0, \nu \sim 0$ (including resummation^[2])

Lam-Tung sum rule: $1 - \lambda = 2\nu$

Experimental data ^[1]: $\nu \sim 30\%$ - Can be interpreted by ISI

^[1] J.S.Conway et al., Phys. Rev. D39 (1989) 92.

^[2] D. Boer et al., Phys. Rev. D74 (2006) 014004.

ν involves transverse spin effects at leading twist ^[2]

If unpolarised DY σ is kept differential on k_T ,

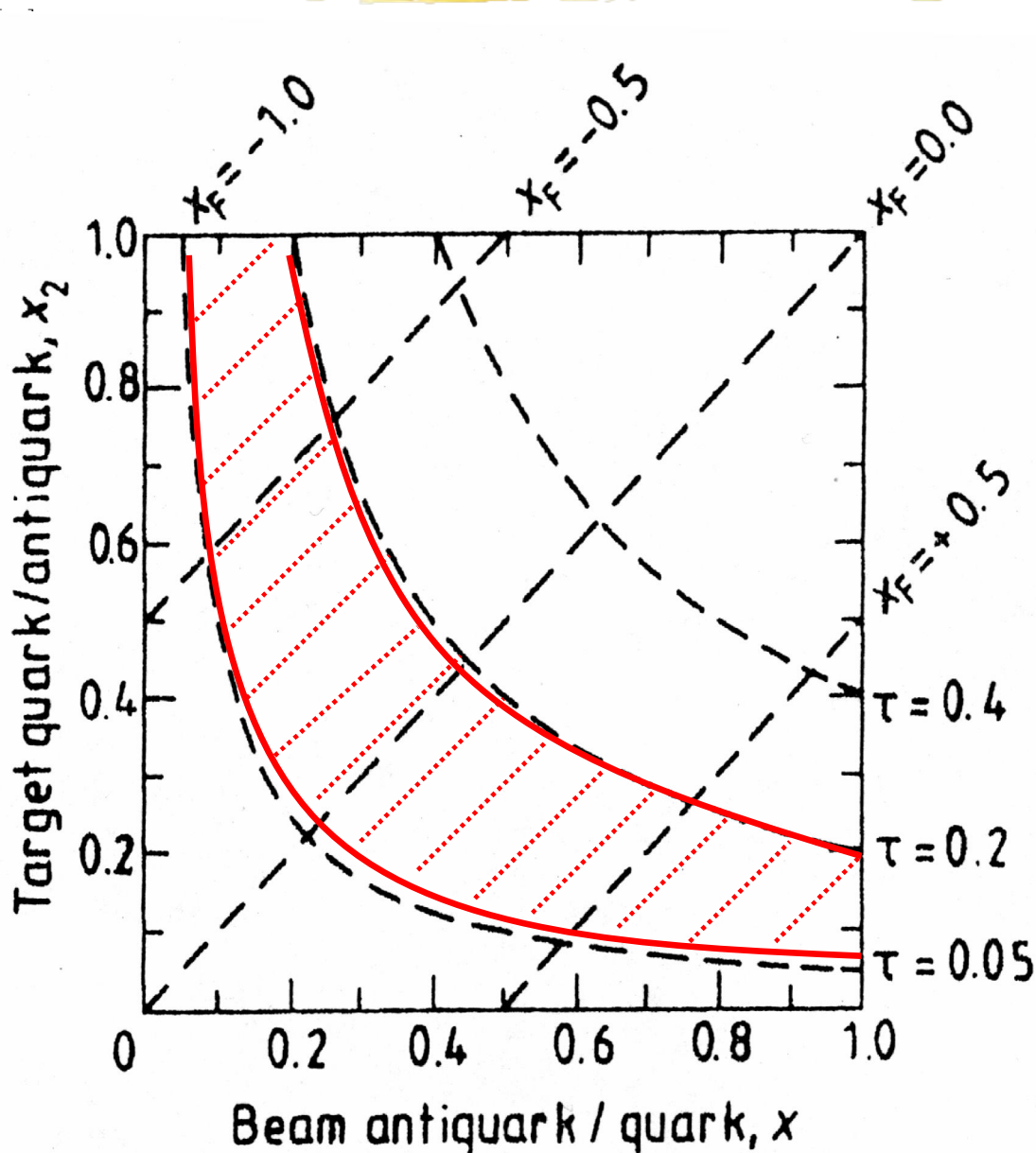
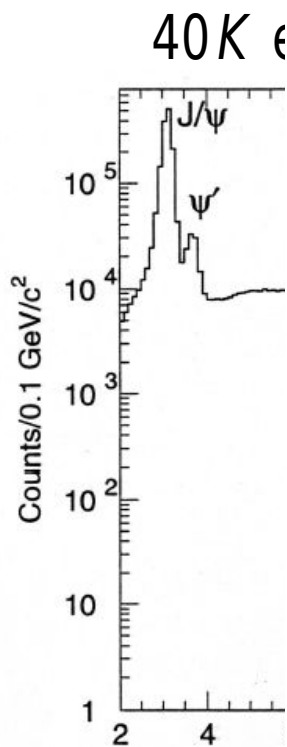
$\cos 2\varphi$ contribution to angular distribution provide:

$$h_1^{\zeta}(\mathbf{x}_2, \kappa_{\zeta}^2) \times \bar{h}_1^{\zeta}(\mathbf{x}_1, \kappa_{\zeta}^2)$$

$\zeta \zeta$

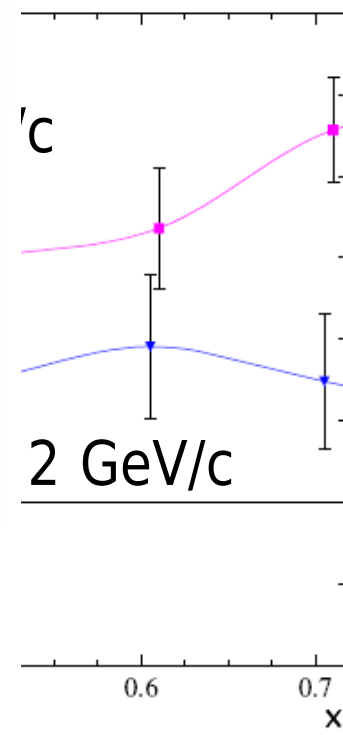
^[2] D. Boer et al., Phys. Rev. D60 (1999) 014012.

Unpolarised Drell-Yan Asymmetries — $\bar{p} p \rightarrow \mu^+ \mu^- X$



GeV/c²

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tion



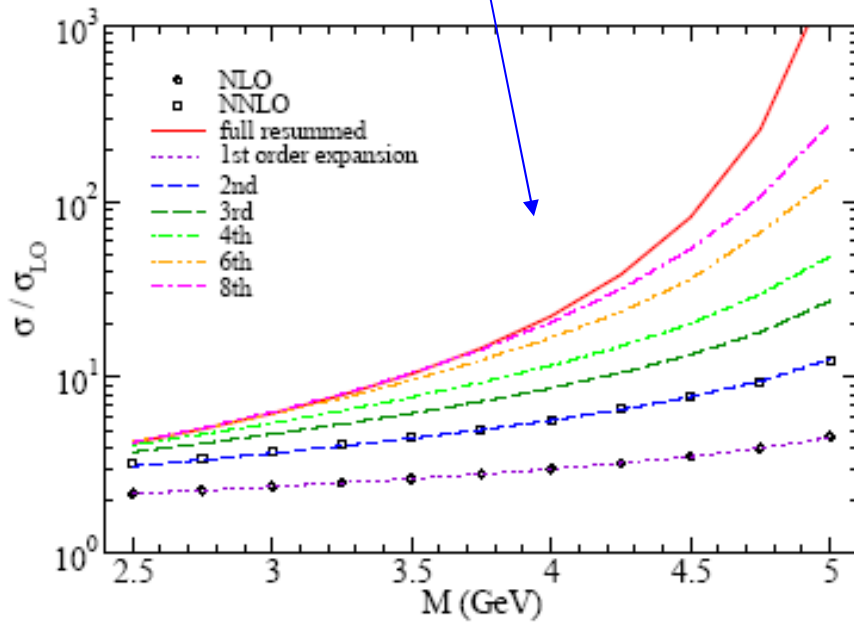
0.2 <

error bars all

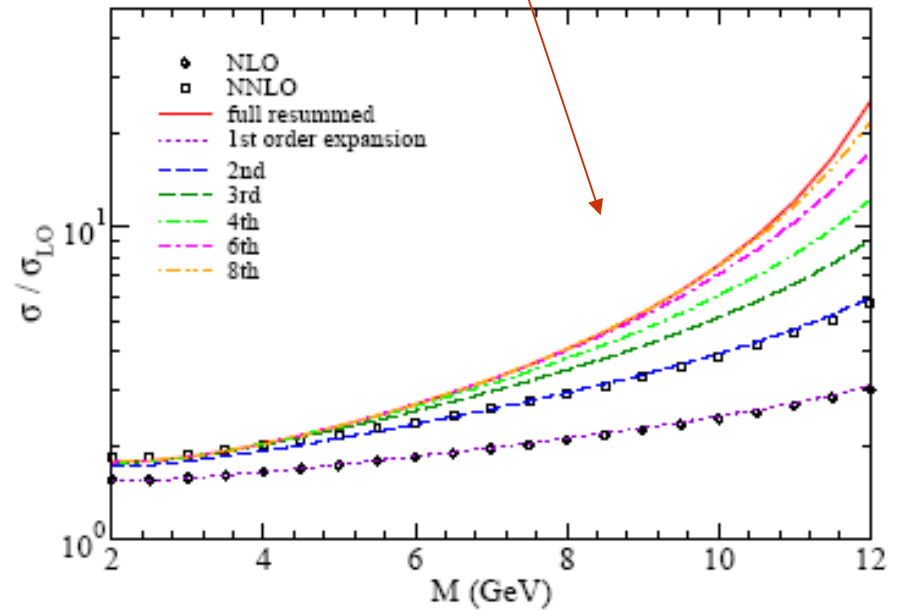
- small asymm
- their depend

Drell-Yan Asymmetries — $\bar{p}p \rightarrow \mu^+ \mu^- X$

$s=30 \text{ GeV}^2$



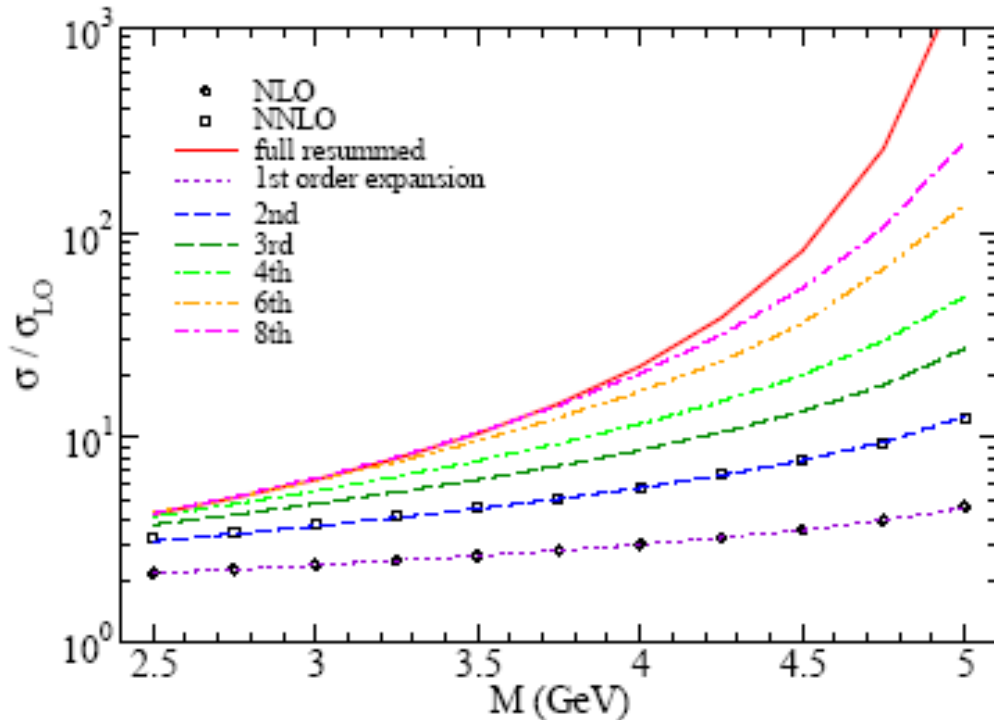
$s=200 \text{ GeV}^2$



At higher energy ($s \sim 200 \text{ GeV}^2$)
perturbative corrections^[1] are sensibly smaller
in the safe region

^[1]H. Shimizu et al., Phys. Rev. D71 (2005) 114007

Unpolarised Drell-Yan — $\bar{p}p \rightarrow \mu^+ \mu^- X$



$$s=30 \text{ GeV}^2$$

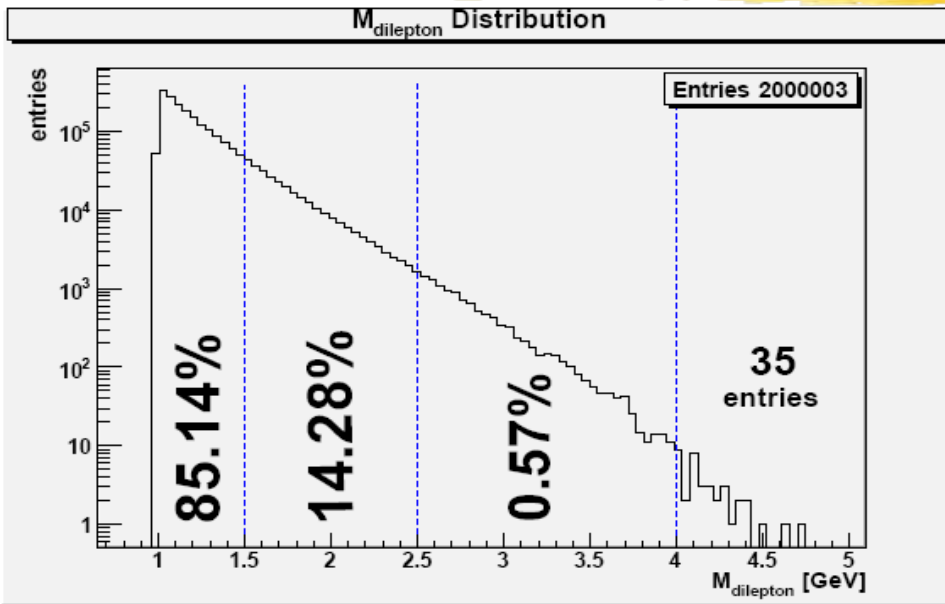
Perturbative corrections^[1] are expected to be large in the PANDA energy range

Unpolarised DY cross-section allow the investigation of:

- limits of the factorisation and perturbative approach
- relation of perturbative and not perturbative dynamics in hadron scattering

^[1]H. Shimizu et al., Phys. Rev. D71 (2005) 114007

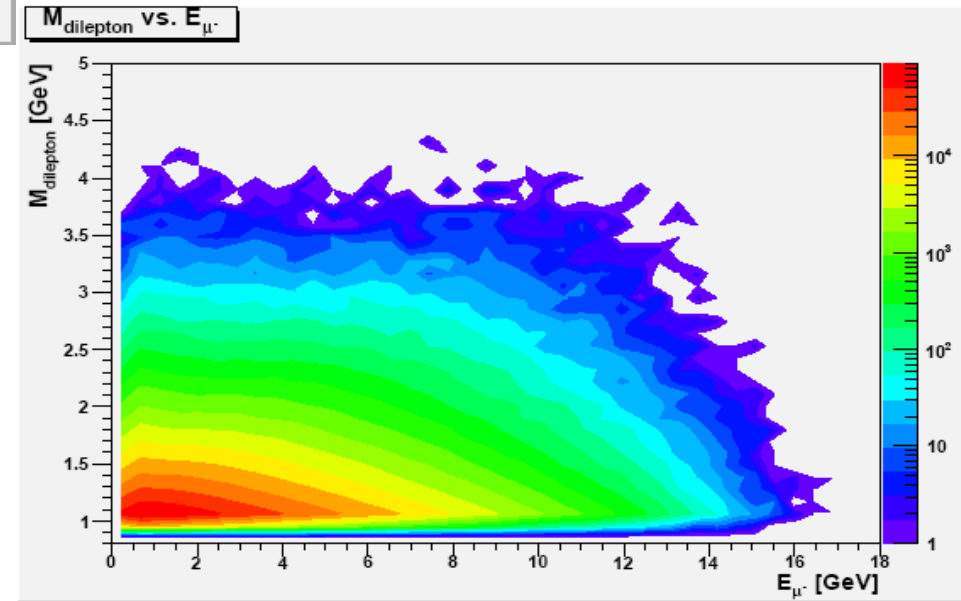
Benchmark channel: DY @ 14 GeV/c — $\bar{p}p \rightarrow \mu^+ \mu^- X$



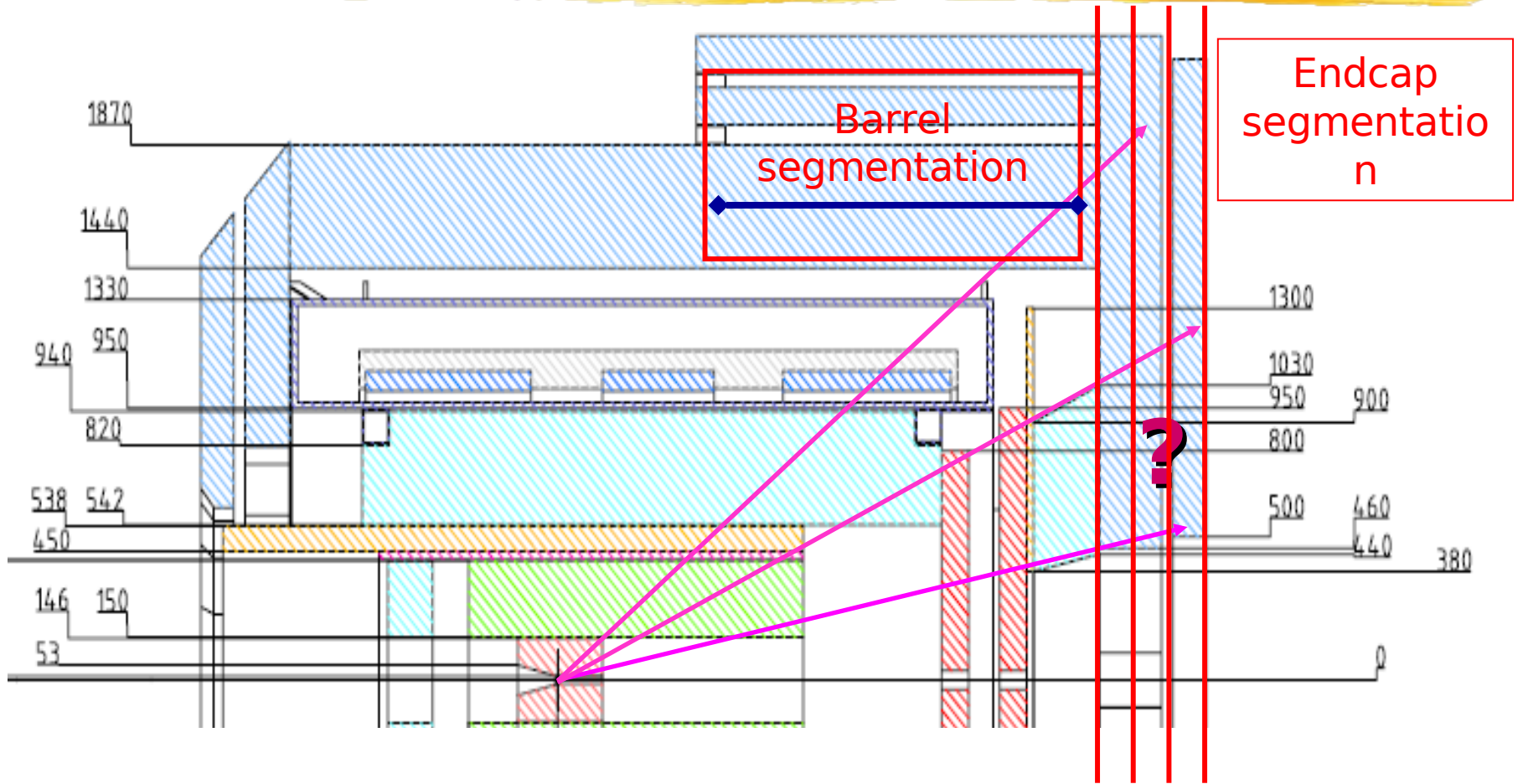
ABianconi Drell - Yan Generator

Focus on two M_{dilepton} ranges

Low Mass **1.1 < M_d < 1.5 GeV**
Medium Mass **1.5 < M_d < 2.5 GeV**



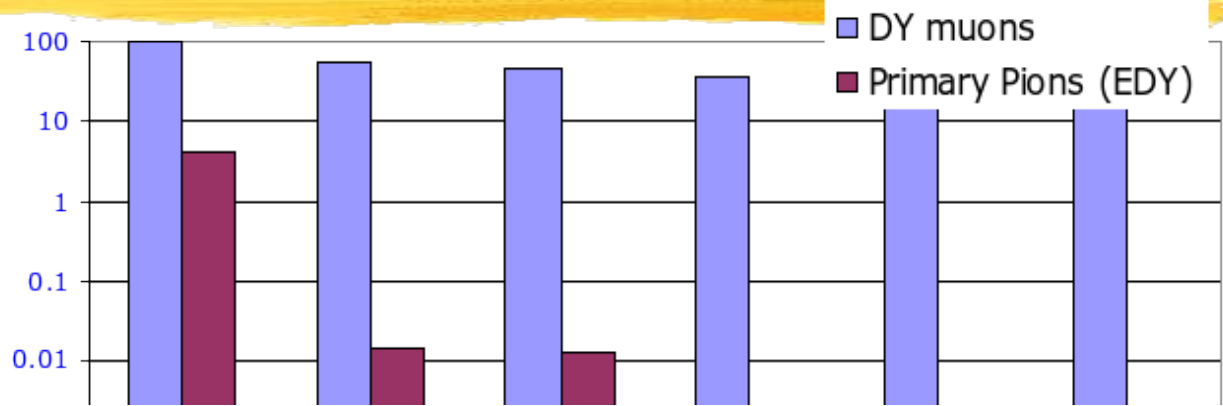
Barrel and endcap segmentation



Background: dipions with the dimuon kinematics

Barrel rejection power

rejection power
after 20 cm iron filter
 $4 \cdot 10^3$

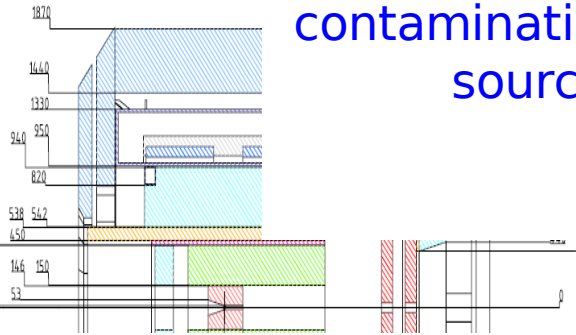


Additional
contamination
sources

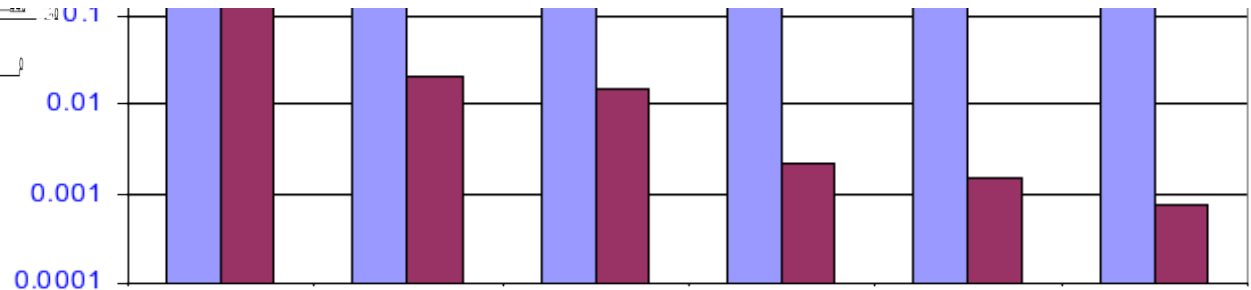
Muons from π decay/primary μ
 10^{-2}

Pions from π interaction/primary π

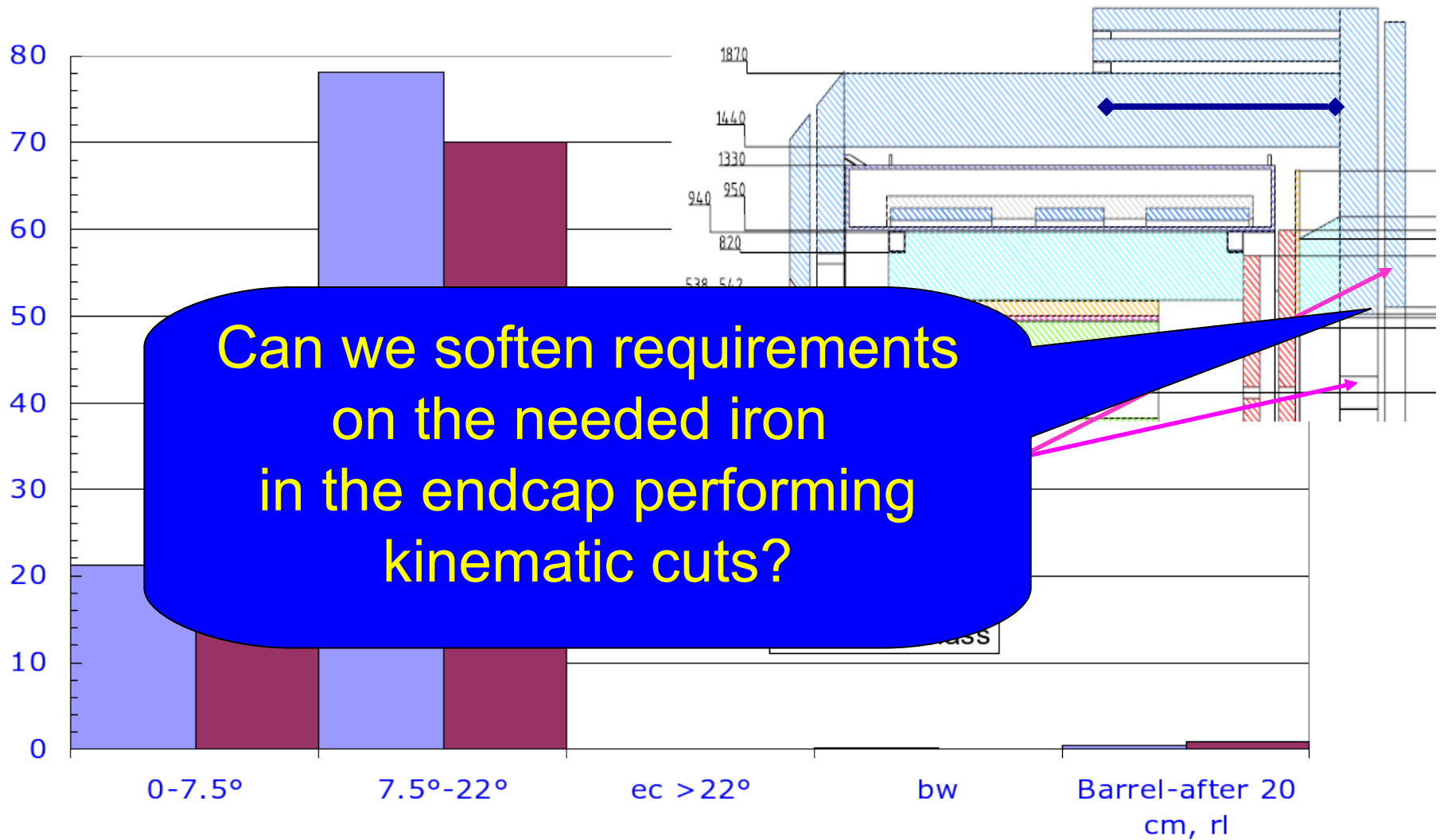
4



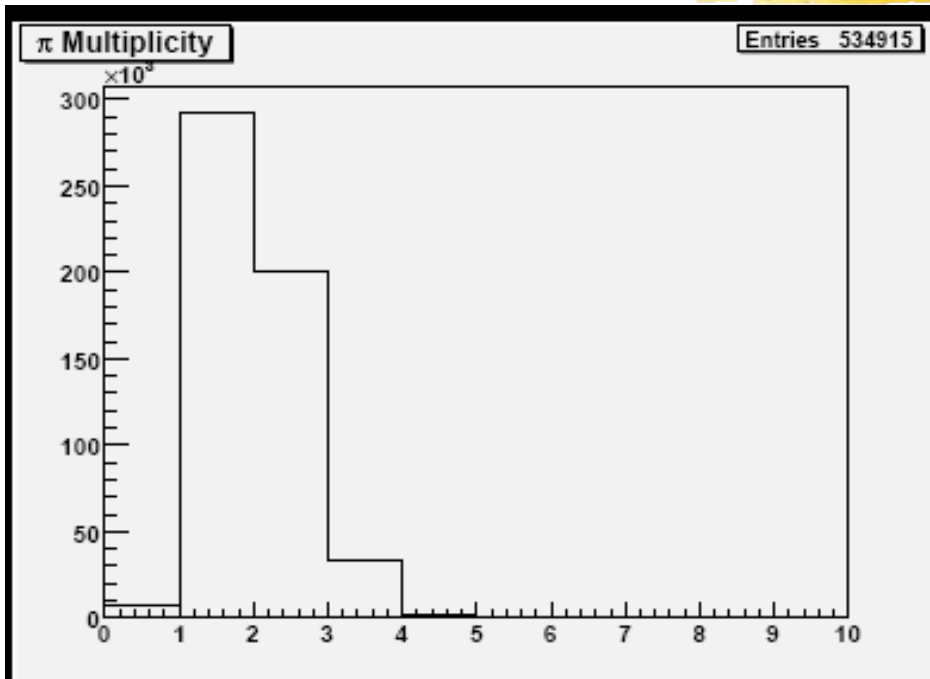
FL= full length
RL=reduced length



Muons partners of those surviving the barrel

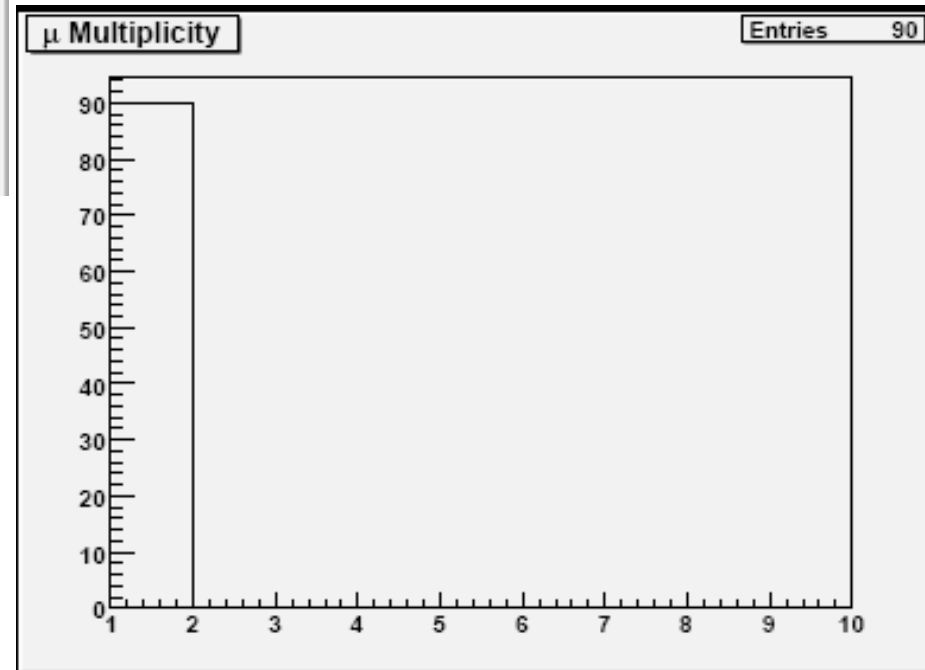


Background simulation with PYTHIA

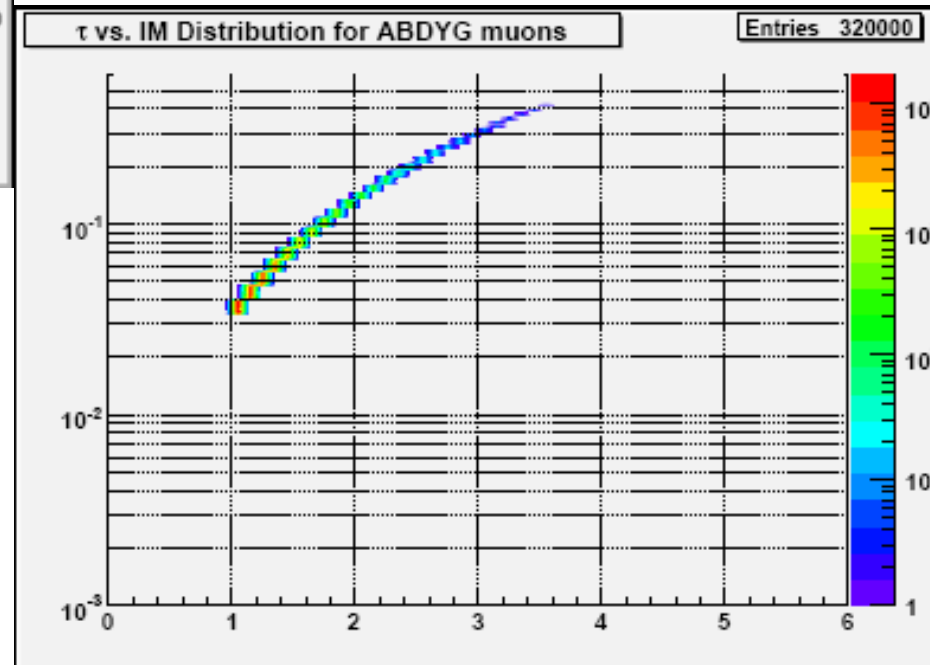
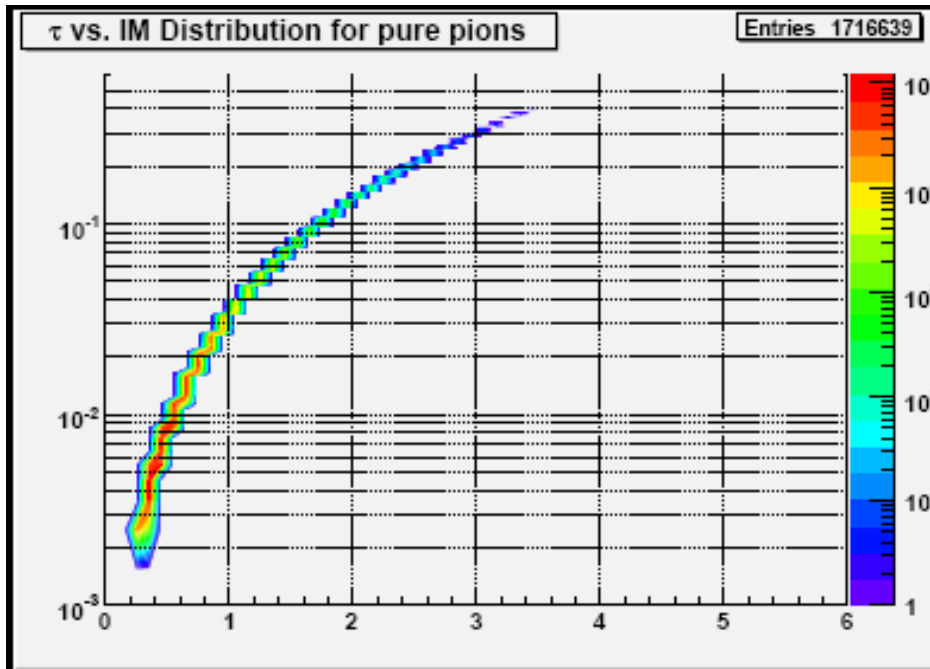


**Most (90%) of the events
have at most two $\pi^+\pi^-$
couples**

**Very few show (at most)
a $\mu^+\mu^-$ couple**

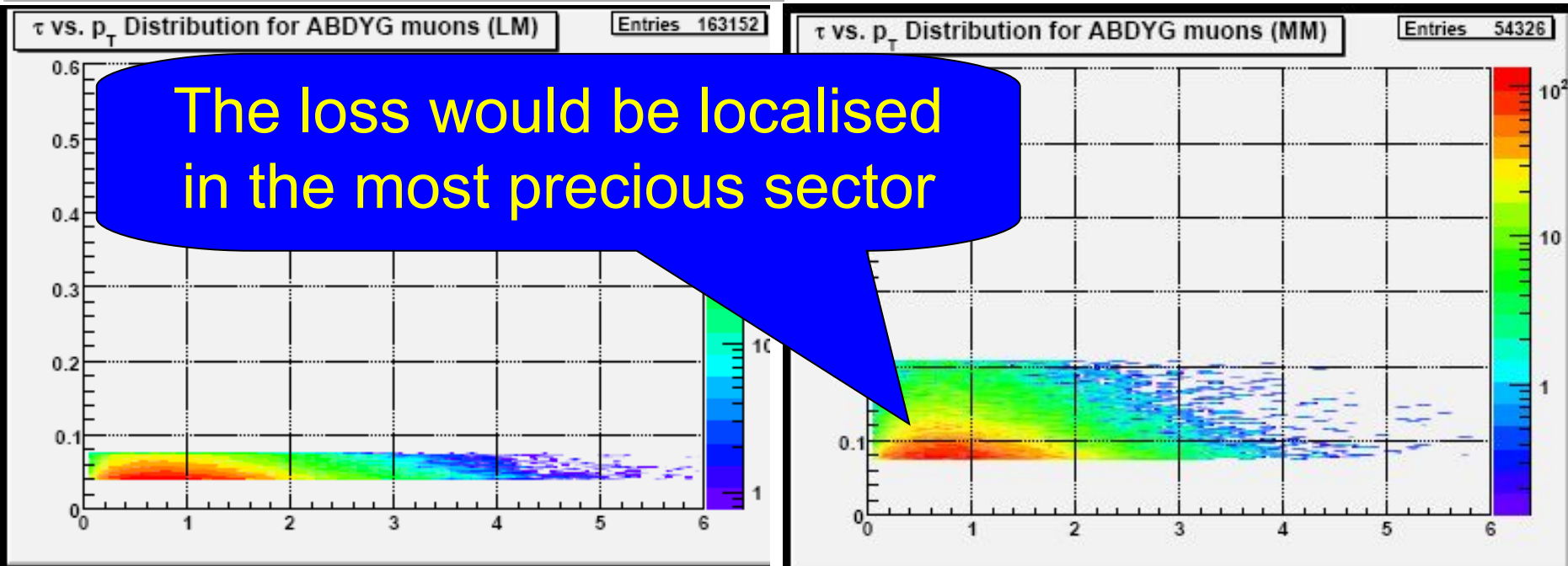
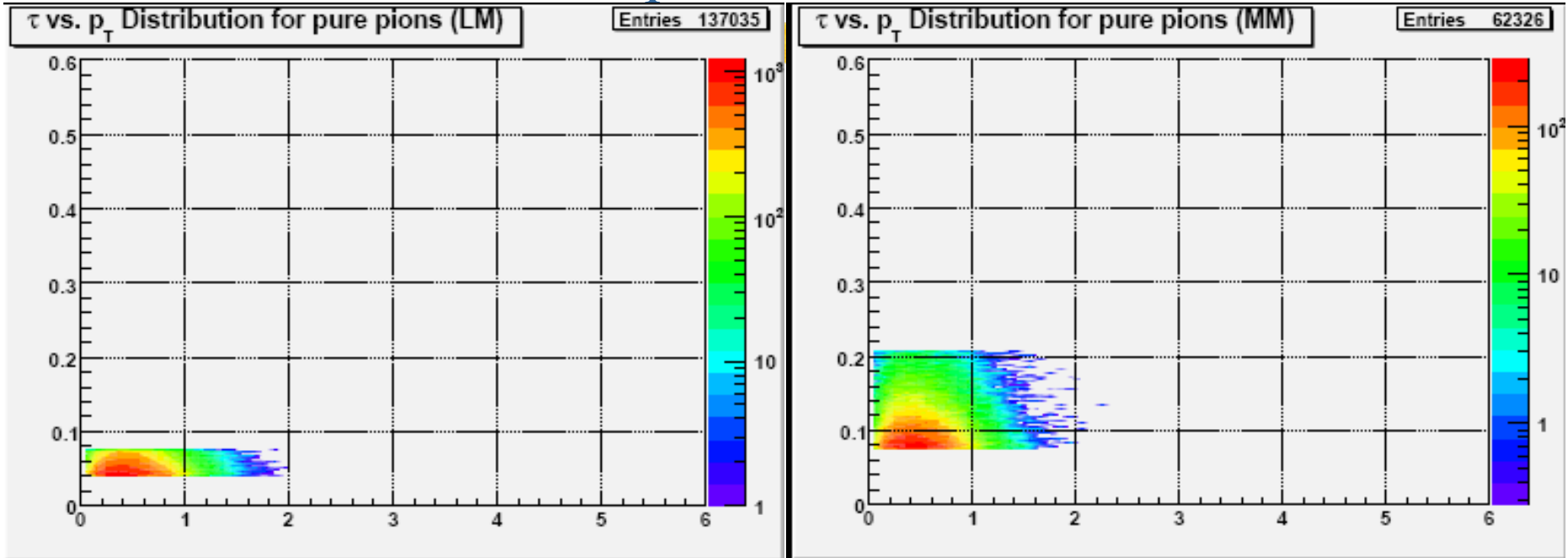


Background and signal kinematics



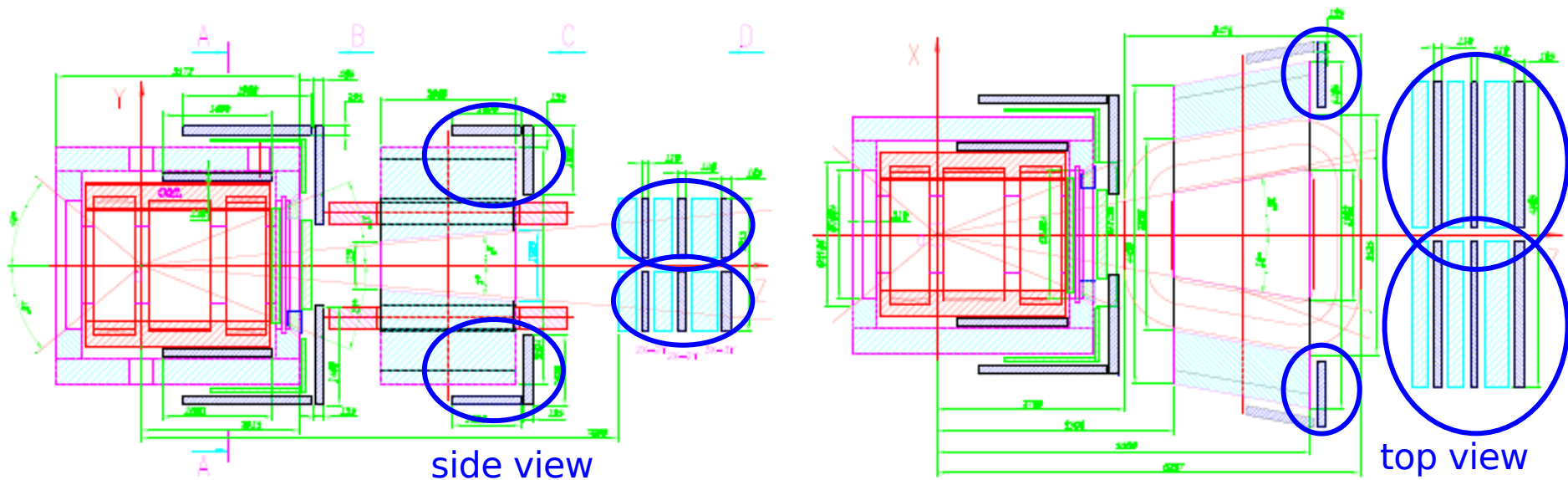
In both cases statistics accumulates in the low IM region

Kinematical cuts are problematic due to statistics loss



Work in progress

- complete background studies simulation
- considering a revised geometry with an enlarged endcap (more Fe-equivalent)
- figuring out requirements in the dipole sector?



Spin physics @ PANDA?

UNPOLARISED DY

$$\cos(2\phi) \text{ asymmetry}^{[1]} \Rightarrow \mathbf{h}_1^{\zeta\zeta}(\mathbf{x}_2, \kappa_{\zeta}^2) \times \bar{\mathbf{h}}_1^{\zeta\zeta}(\mathbf{x}_1, \kappa_{\zeta}^2)$$

Unpolarised cross section

^[2] D. Boer et al., Phys. Rev. D60 (1999) 014012.

ANTIPROTONS!!

DY azimuthal asymmetries
not suppressed by nonvalence-like contributions.

DETECTOR STUDIES IN PROGRESS

Question time



THANK YOU!

Drell-Yan Asymmetries — $\bar{p}^\uparrow p^\uparrow \rightarrow \mu^+ \mu^- X$

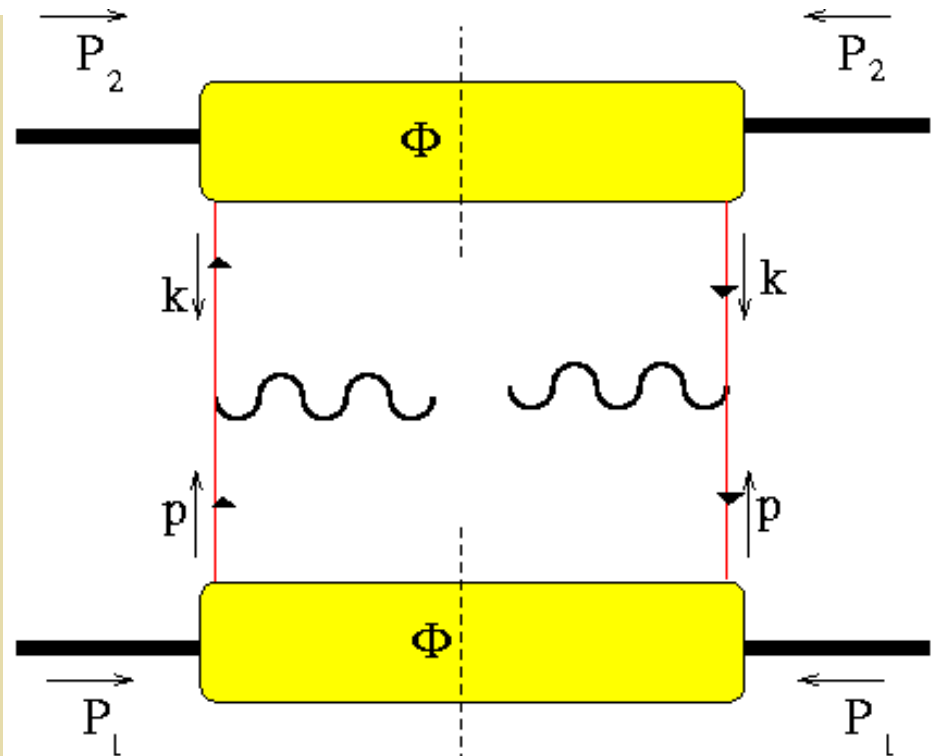
Uncorrelated quark helicities \Rightarrow access chirally-odd functions



TRANSVERSITY

Ideal because:

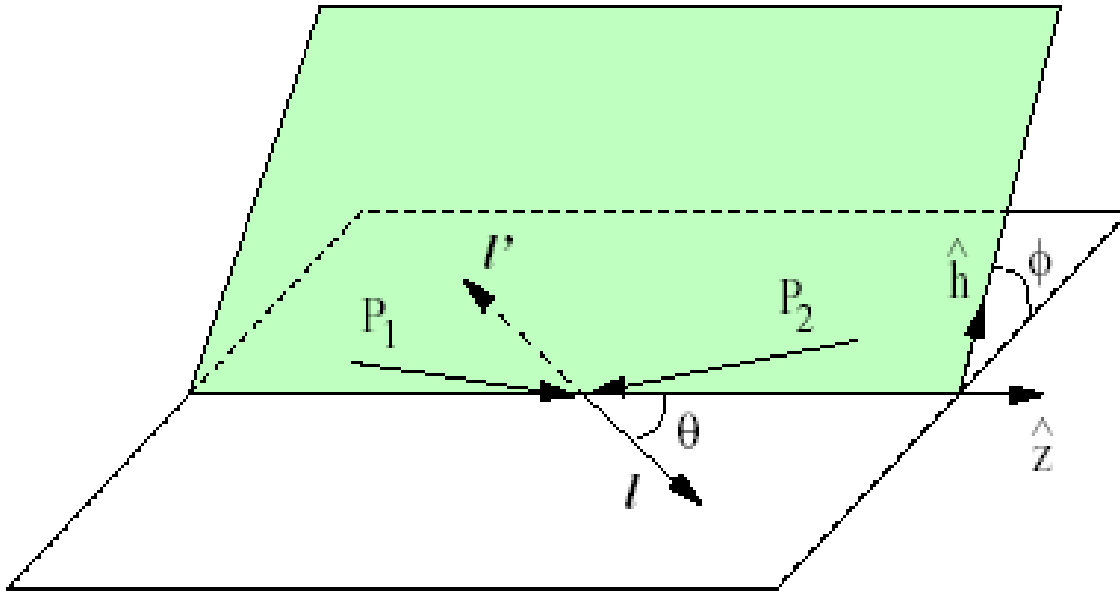
- h_1 not to be unfolded with fragmentation functions
- chirally odd functions not suppressed (like in DIS)



Drell-Yan Asymmetries — $\bar{p}^\uparrow p^\uparrow \rightarrow \mu^+ \mu^- X$

$$A_{LL} = \frac{\sum_a e_a^2 g_1^a(x_1) g_1^{\bar{a}}(x_2)}{\sum_a e_a^2 f_1^a(x_1) f_1^{\bar{a}}(x_2)} \quad A_{TT} = \frac{\sin^2 \theta \cos 2\phi}{1 + \cos^2 \theta} \frac{\sum_a e_a^2 h_1^a(x_1) h_1^{\bar{a}}(x_2)}{\sum_a e_a^2 f_1^a(x_1) f_1^{\bar{a}}(x_2)}$$

$$A_{LT} = \frac{2 \sin 2\theta \cos \phi}{1 + \cos^2 \theta} \frac{M}{\sqrt{Q^2}} \frac{\sum_a e_a^2 \left(g_1^a(x_1) x_2 g_T^{\bar{a}}(x_2) - x_1 h_L^a(x_1) h_1^{\bar{a}}(x_2) \right)}{\sum_a e_a^2 f_1^a(x_1) f_1^{\bar{a}}(x_2)}$$



lepton plane (cm)

Collins-Soper frame: ^[1]Phys. Rev. D16 (1977) 2219.

To be corrected for:

$$\frac{1}{P_{\bar{p}} f P_{\bar{p}}}$$

Drell-Yan Asymmetries — $\bar{p} p^\uparrow \rightarrow \mu^+ \mu^- X$

$$\frac{1}{\sigma} \frac{d\sigma}{d\Omega} \propto \left(1 + \cos^2 \theta + \frac{\nu}{2} \sin^2 \theta \cos 2\varphi + \rho |S_{1T}| \sin^2 \theta \sin(\varphi - \varphi_{S_1}) + \dots \right)$$

$$\lambda \sim 1, \mu \sim 0$$

$$A_T = \frac{|S_{1T}| \nu \sin 2\theta \sin(\varphi - \varphi_{S_1})}{1 + \cos^2 \theta} \frac{N}{\sqrt{D}}$$

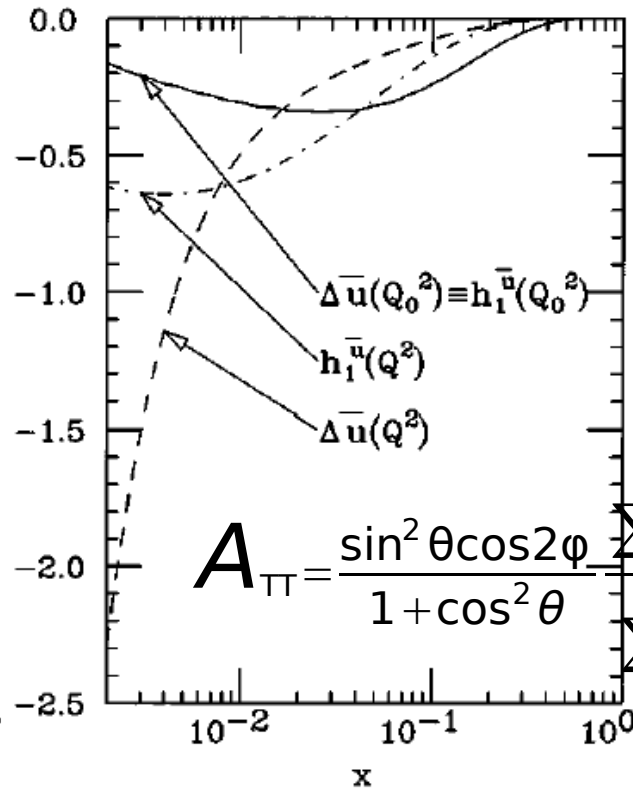
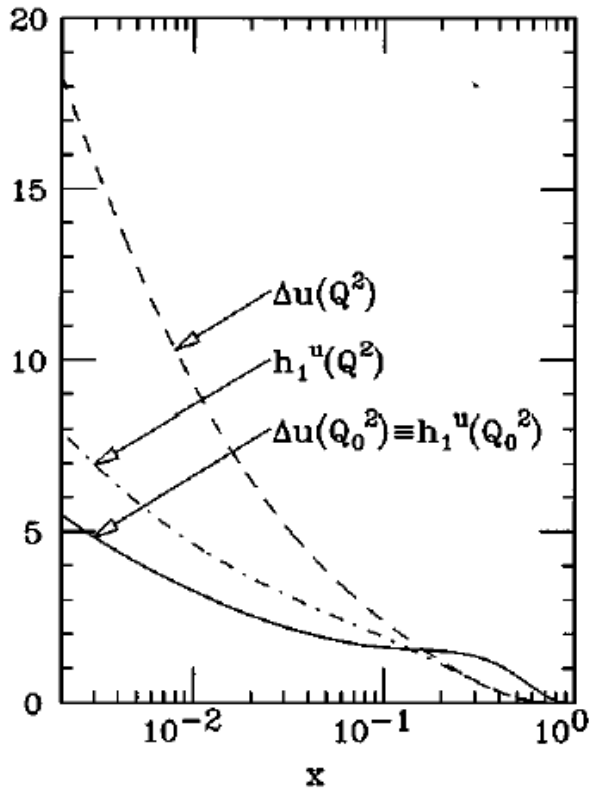
Even unpolarised \bar{p} beam on polarised p,
 or polarised \bar{p} on unpolarised p
 are powerful tools
 to investigate κ_T dependence of QDF

Drell-Yan Asymmetries — $\bar{p}^\uparrow p^\uparrow \rightarrow \mu^+ \mu^- X$

RHIC energies: $\sqrt{s}=100$ GeV $M^2=100 \Rightarrow \tau \leq 10^{-2} \Rightarrow$ small x_1 and/or x_2

$h_1^a(x, Q^2)$ evolution much slower^[1] than $\Delta q(x, Q^2)$ and $q(x, Q^2)$ at small x

A_{TT} @ RHIC very small, smaller \sqrt{s} would help^[1]



$$Q_0^2 = 0.23 \text{ GeV}^2$$

$$Q_a^2 = 25 \text{ GeV}^2$$

$$A_{TT} = \frac{\sin^2 \theta \cos 2\phi}{1 + \cos^2 \theta} \frac{\sum_a e_a^2 h_1^a(x_1) h_1^a(x_2)}{\sum_a e_a^2 f_1^a(x_1) f_1^a(x_2)}$$

^[1]Barone, Colarco and Drago, Phys.Rev. D56 (1997) 527.

Drell-Yan Asymmetries — $\bar{p}^\uparrow p^\uparrow \rightarrow \mu^+ \mu^- X$

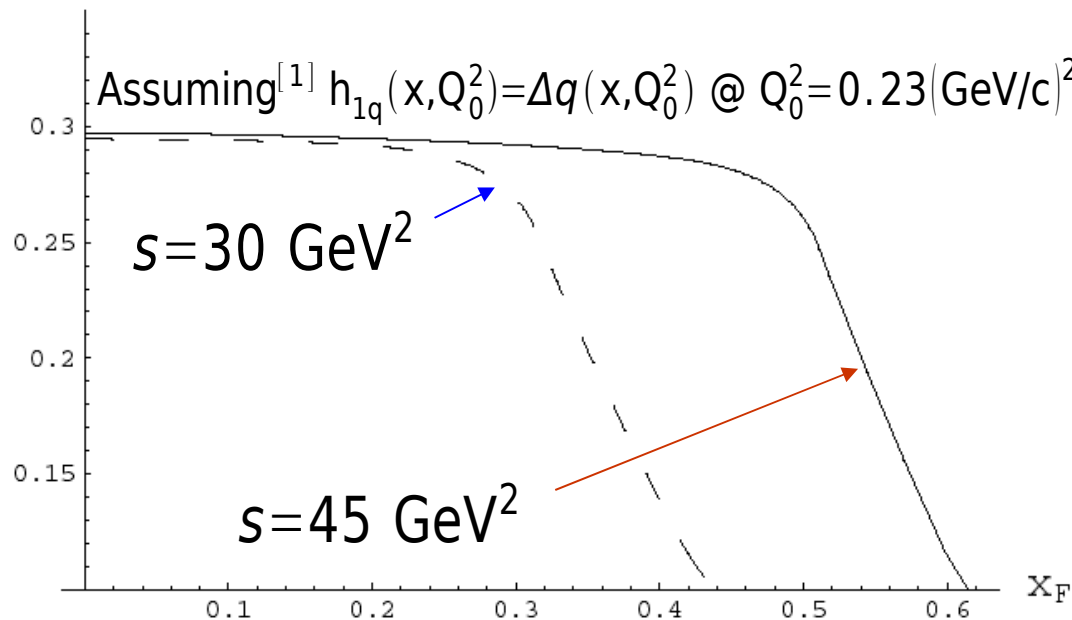
$$A_{\text{TT}} = \hat{a}_{\text{TT}} \frac{\sum_q e_q^2 [h_{1q}^{\bar{p}}(x_1) h_{1\bar{q}}^p(x_2) + h_{1\bar{q}}^{\bar{p}}(x_1) h_{1q}^p(x_2)]}{\sum_q e_q^2 [q^{\bar{p}}(x_1) \bar{q}^p(x_2) + \bar{q}^{\bar{p}}(x_1) q^p(x_2)]}$$

$$\stackrel{\substack{i \\ \text{large } x}}{\sim} \hat{a}_{\text{TT}} \frac{\sum_q e_q^2 h_{1q}^p(x_1) h_{1q}^p(x_2)}{\sum_q e_q^2 q^p(x_1) q^p(x_2)}$$

A_{TT} still small @ large \sqrt{s} and M^2 due to slow evolution of $h_1^a(x, Q^2)$

Large A_{TT} expected^[1] for \sqrt{s} and M^2 not too large and τ not too small

$A_{\text{TT}}^{\bar{p}^\uparrow p^\uparrow} / a_{\text{TT}}$



HESR: $s_{\text{max}} = 30 \div 45 \text{ GeV}^2$
 $M^2 \geq M_{J/\psi}^2 \rightarrow \tau \geq 0.3$

A_{TT} direct access
to valence quark h_1



$h_{1q_V}(x_1) \times h_{1q_V}(x_2)$

^[1]M. Anselmino et al., Phys. Lett. B594 (2004) 97.

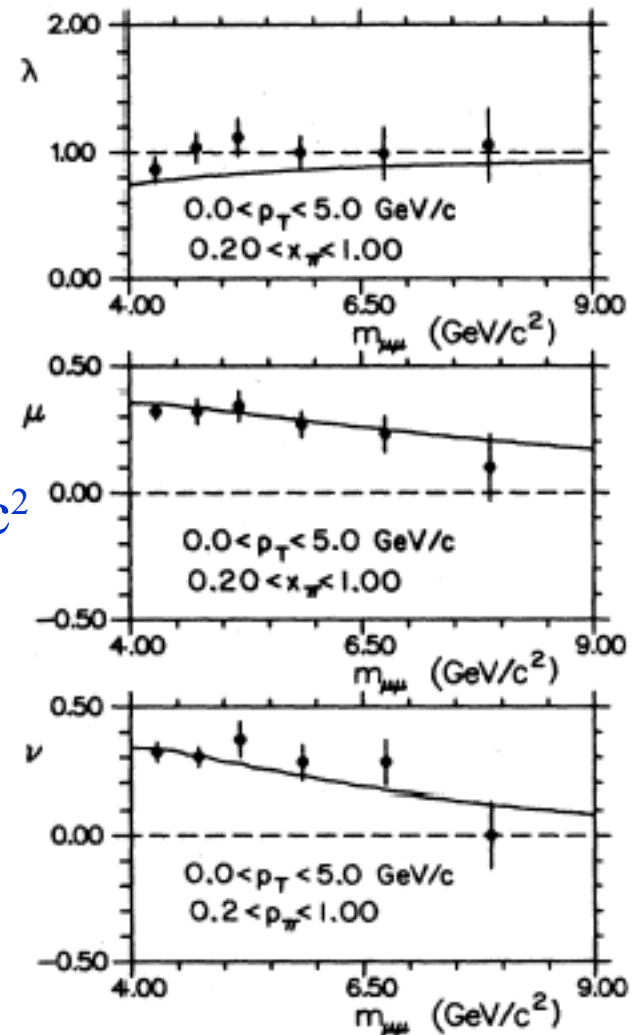
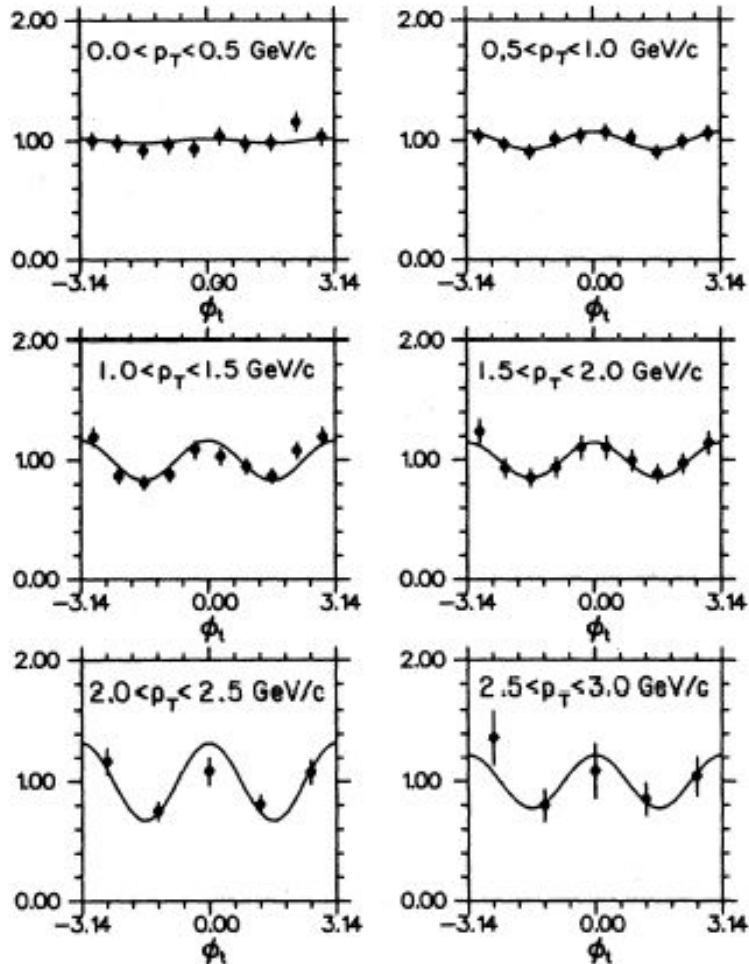
Angular distribution in CS frame

E615 @ Fermilab

$\pi-N \rightarrow \mu+\mu-X$ @ 252 GeV/c

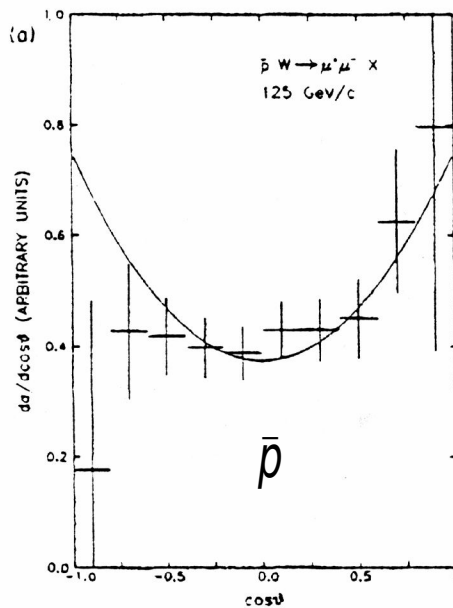
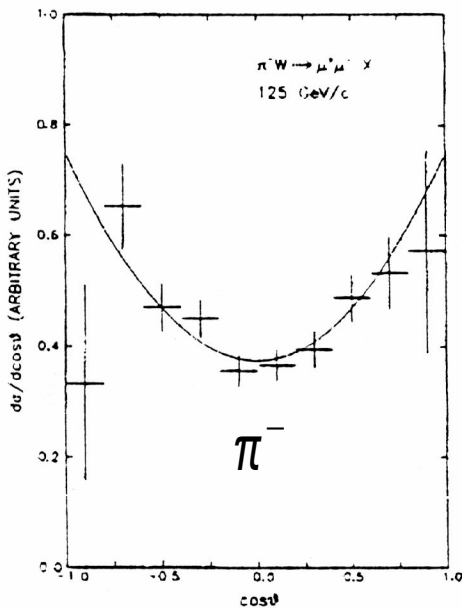
$-0.6 < \cos\vartheta < 0.6$

$4 < M < 8.5$ GeV/c²



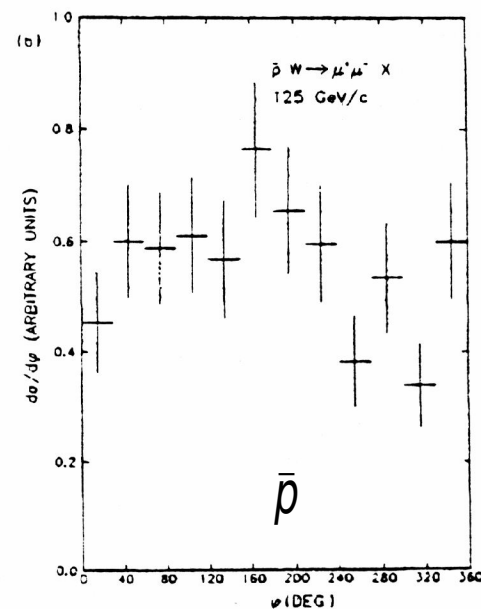
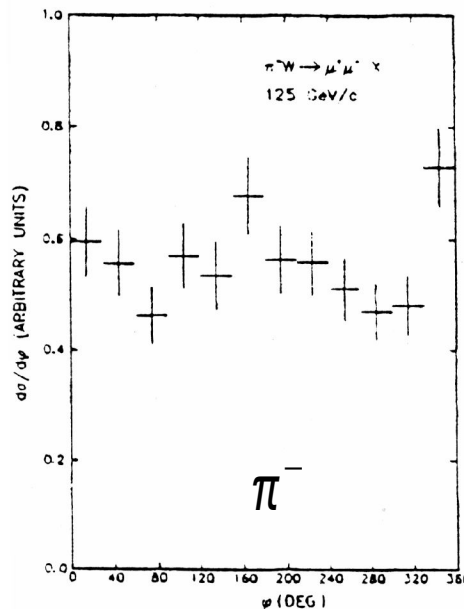
- cut on P_T selects asymmetry
- 30% asymmetry observed for π

Angular distributions for \bar{p} and π^- - π -N, \bar{p} N @ 125 GeV/c



- $\frac{d\sigma}{d\cos\theta}$ vs $\cos\theta$

- $\frac{d\sigma}{d\phi}$ vs ϕ



Transverse Single Spin Asymmetries

$$A_N = \frac{d\sigma^\uparrow - d\sigma^\downarrow}{d\sigma}$$

E704 Tevatron FNAL 200 GeV/c

$p \rightarrow \pi X$

- π Production @ 100 GeV

π^+ : $A_N > 0$; π^- : $A_N < 0$

- A_N similar for \sqrt{s}

A_N related to funda

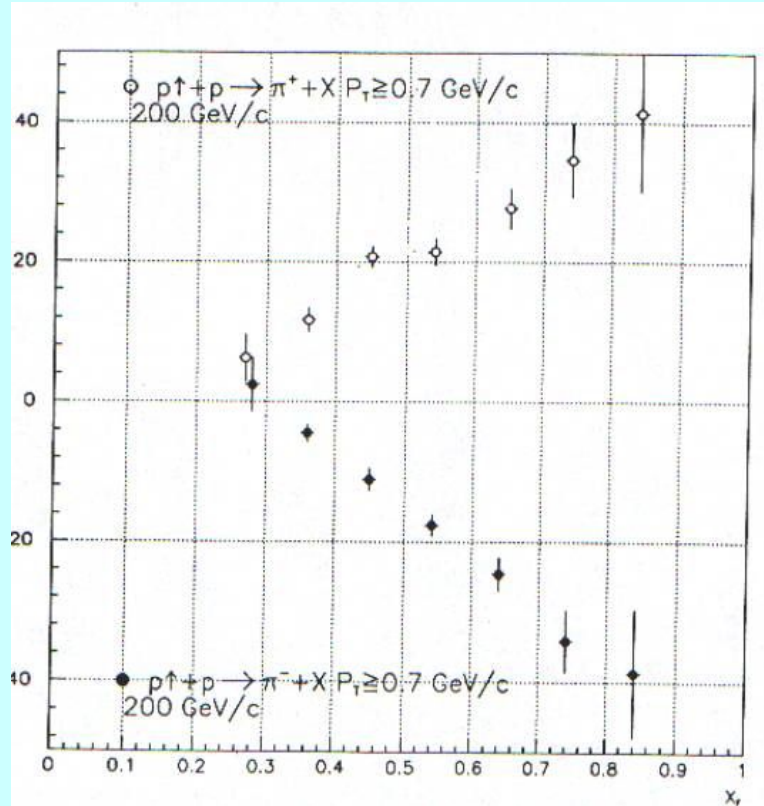
New experiment with \sqrt{s} region:

- new data available

- $A_{N, \{\bar{p}p^\uparrow \rightarrow \pi X\}}$ vs $A_{N, \{\bar{p}^\uparrow p \rightarrow \pi X\}}$

- DY-SSA (A_T) possible only @ RHIC, $p^\uparrow p$ -scattering:

$$\sigma_{\bar{p}p}^{DY} @ \text{smaller } s \gg \sigma_{pp}^{DY} @ \text{large } s$$



polarisation

orientation

kinematical